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Monte Carlo Investigation of Refractive Index Gas Thermometry Uncertainties

ABSTRACT

Refractive Index Gas Thermometry (RIGT) is a primary thermometry technique whereby the thermodynamic temperature of a working gas is determined via measurements of the refractive index of the gas. Monte Carlo calculations are presented for a microwave resonance-based RIGT arrangement under development at NRC, employing a quasi-spherical copper resonator and helium as the working gas.

Results indicate that RIGT could be a useful primary thermometry technique to supplement Acoustic Gas Thermometry (AGT) and Dielectric Constant Gas Thermometry (DCGT) thermodynamic temperature measurements between 25 K and 150 K.



Figure 1. Quasi-spherical microwave resonator & gas pressure vessel
From Rourke (2014) [1]

INTRODUCTION

- Consultative Committee for Thermometry Working Group 4 (WG4) published consensus estimate of ITS-90 temperature deviation from thermodynamic temperature – Fischer (2011) [2]
- 24.5 K – 77 K region: special interest
 - Gap in AGT results of Pitre (2006) [3]
 - Disagreement between AGT results of Pitre (2006) [3] and DCGT results of Gaiser (2010) [4]
- RIGT: similarities to both AGT & DCGT ⇒ tie-breaker?
- RIGT uncertainties: low enough to contribute meaningfully to $T - T_{90}$?
 - Treated at 1 standard deviation level in present work

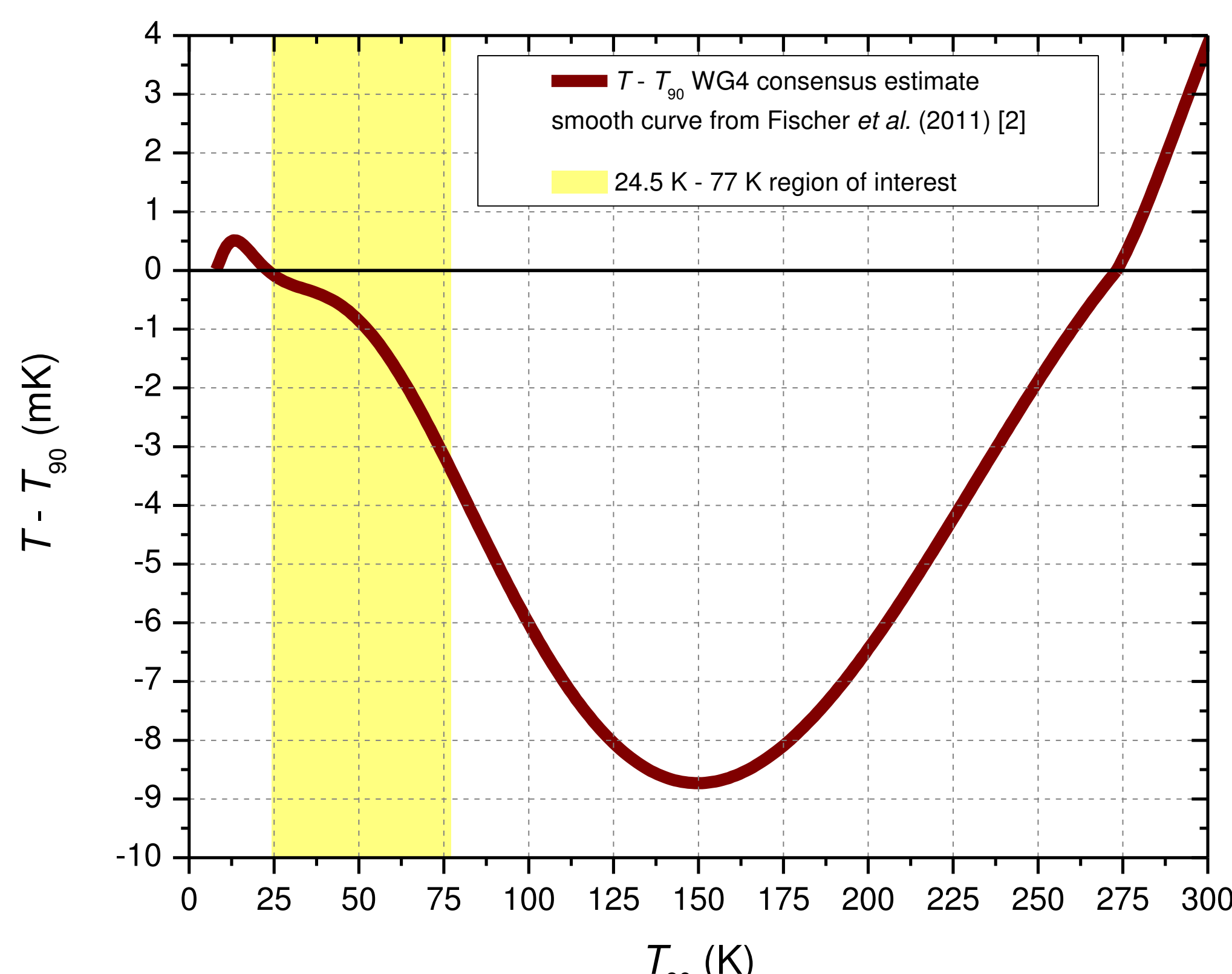


Figure 2. Consensus estimate of ITS-90 temperature T_{90} deviation from thermodynamic temperature T
From Fischer (2011) [2]

MAIN EQUATIONS

- Gas refractive index n_{expt} obtained experimentally, in which c_0 is the speed of light in vacuum and $\langle \xi_{\text{corr}} \rangle$ is the average corrected microwave eigenvalue, as per Rourke (2014) [1]

$$n_{\text{expt}}^2 - 1 = \left\{ \frac{c_0}{2\pi \cdot \langle f + g \rangle(T)} \cdot \frac{\langle \xi_{\text{corr}} \rangle}{a_0(T) \cdot [1 - \kappa_T(T) \cdot p]} \right\}^2 - 1$$

- Copper isothermal compressibility κ_T

$$\kappa_T = \frac{1}{B_s} + \frac{9 \cdot \alpha_L^2(T) \cdot T \cdot a_0^3(T)}{\rho_{\text{Cu}}(T = 293 \text{ K}) \cdot c_p(T) \cdot a_0^3(T = 293 \text{ K})}$$

- Gas pressure p virial equation in terms of the molar gas density ρ

$$p \approx RT \cdot [\rho + B_p(T) \cdot \rho^2 + C_p(T) \cdot \rho^3 + D_p(T) \cdot \rho^4]$$

- Gas pressure virial equation inverted by iterative substitution, as per Moldover (1998) [5], to get the molar gas density ρ in terms of the gas pressure p

$$\rho \approx p \left\{ RT + \left[\frac{B_p(T) \cdot p}{1 + \frac{B_p(T) \cdot p}{RT + B_p(T) \cdot p} + \frac{C_p(T) \cdot p^2}{RT(RT + 2B_p(T) \cdot p)}} \right] + \frac{C_p(T) \cdot p^2}{RT + 2B_p(T) \cdot p} + \frac{D_p(T) \cdot p^3}{R^2 T^2} \right\}$$

- Gas relative magnetic permeability μ_r virial equation

$$(\mu_r - 1)/(\mu_r + 2) \approx A_\mu \rho$$

- Gas relative dielectric permittivity ϵ_r virial equation

$$(\epsilon_r - 1)/(\epsilon_r + 2) \approx A_\epsilon \rho + B_\epsilon(T) \cdot \rho^2 + C_\epsilon \rho^3$$

- Gas refractive index n_{calc} obtained theoretically

$$n_{\text{calc}}^2 - 1 = (\mu_r \epsilon_r) - 1$$

PARAMETER VALUES AND UNCERTAINTIES

- Input temperature T** uncertainty (temperature stability & ITS-90 realization) set to 0.2 mK
- Pressure p** uncertainties from NRC's CMCs ($k = 1$)
 - > 5 kPa – 350 kPa: 1.00×10^{-5} relative uncertainty
 - > 350 kPa – 1750 kPa: 1.80×10^{-5} relative uncertainty
 - > 1750 kPa – 7000 kPa: 2.50×10^{-5} relative uncertainty
- Average half-width-corrected microwave frequency $\langle f + g \rangle$** uncertainty and **resonator vacuum radius $a_0(T)$** values & uncertainties from cubic spline interpolation of Rourke (2014) [1] experimental measurements
- OFHC copper material parameters contributing to the **resonator isothermal compressibility κ_T** : adiabatic bulk modulus B_s , linear thermal expansion coefficient α_L , specific heat at constant pressure c_p , and density at 293 K $\rho_{\text{Cu}}(T = 293 \text{ K})$ – values & uncertainties from Simon (1992) [6]
 - > κ_T uncertainty dominated by uncertainty in B_s
- Molar gas constant $R = 8.3144621(75) \text{ J mol}^{-1} \text{ K}^{-1}$ from CODATA (2010)
- Helium gas density virial coefficients $B_p(T)$, $C_p(T)$, $D_p(T)$ values & uncertainties from cubic spline interpolation of Shaul (2012) [7]
- Helium gas diamagnetic virial coefficient $A_\mu = 4\pi\chi_0/3 = -0.000007921(4) \text{ cm}^3 \text{ mol}^{-1}$ from Bruch (2000) [8], following treatment of Moldover (2014) [9]
- Helium gas molar polarizability in the limit of zero density $A_\epsilon = 0.51725419(10) \text{ cm}^3 \text{ mol}^{-1}$ from Lach (2004) [10], following treatment of Moldover (2014) [9]
- Helium gas dielectric virial coefficient $B_\epsilon(T)$ from 5th-order polynomial fit Rizzo (2002) [11] calculated table values; uncertainties are quadrature combinations of 3 sources
 - > Rizzo's expectation of convergence to within 0.1%
 - > Digital display uncertainty for Rizzo table values printed to 4 decimal places without printed uncertainties
 - > $5.4 \times 10^{-5} \text{ cm}^6 \text{ mol}^{-2}$ standard deviation of 5th-order polynomial fit residuals
- Helium gas dielectric virial coefficient $C_\epsilon = -0.6(4) \text{ cm}^9 \text{ mol}^{-3}$** from averaging several experimental [12-15] and calculated [16] values at different temperatures; uncertainty from standard deviation of these values
 - > Roughly T -independent, at least within published uncertainties

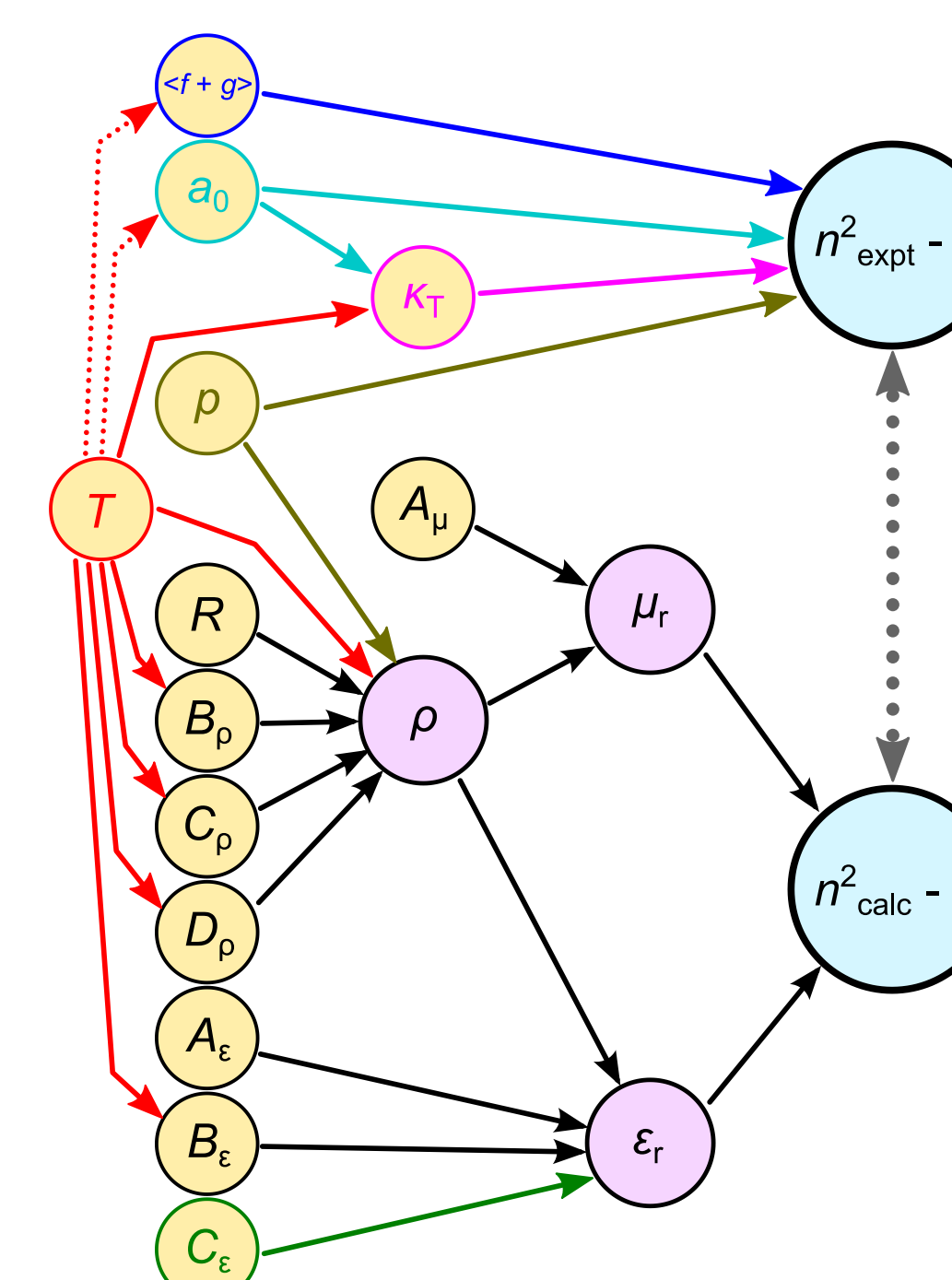


Figure 3. Relational diagram of RIGT equation parameters
Parameters with yellow backgrounds are sampled via Monte Carlo trials

MONTE CARLO CALCULATION DETAILS

- For each input parameter at a given $(T, p, \text{microwave mode})$ point:
 - Input parameter Gaussian distribution sampled 100,000 times
 - Resulting $n^2 - 1$ standard deviation translated to a thermodynamic temperature T_{therm} distribution to get δT_{therm} , using relationship between T distribution width and $n^2_{\text{calc}} - 1$ distribution width

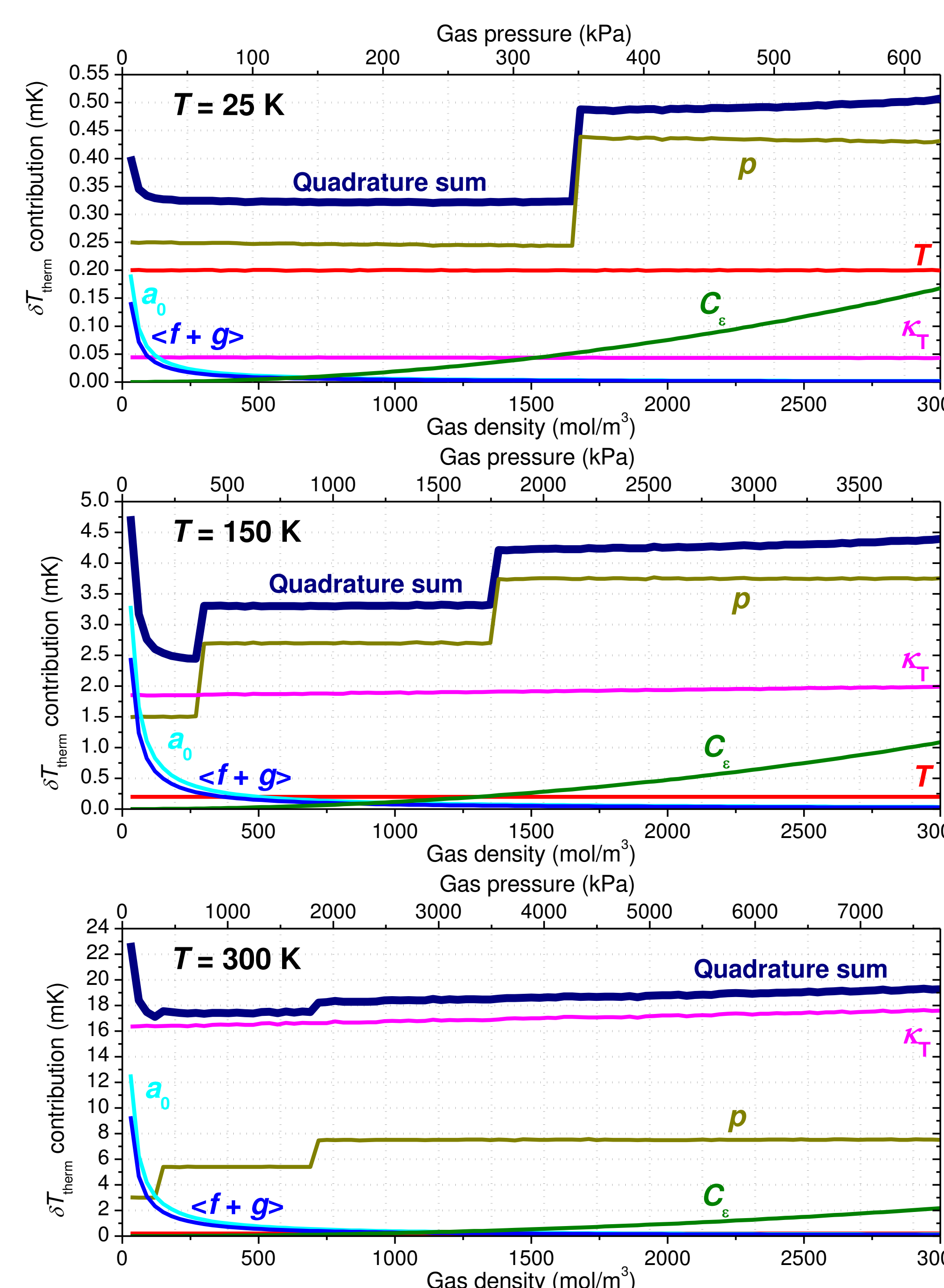


Figure 4. Dominant Monte Carlo contributions to RIGT thermodynamic temperature uncertainty
Calculated at various temperatures and gas pressures, for TM11 microwave mode

RESULTS

- RIGT thermodynamic temperature uncertainties dominated by uncertainties in gas pressure p and, at higher temperatures, resonator compressibility κ_T
- Best overall uncertainties obtained when $p \leq 350 \text{ kPa}$
- Results for microwave modes TE11, TM12, TE12 and TM13 (not shown) similar to TM11, with increased uncertainty contributions from a_0 and $\langle f + g \rangle$

Temperature T	Gas pressure p	Gas density ρ	δT_{therm} quadrature sum
25 K	350 kPa	1700 mol/m ³	0.32 mK
50 K	350 kPa	840 mol/m ³	0.57 mK
75 K	350 kPa	560 mol/m ³	0.88 mK
150 K	350 kPa	280 mol/m ³	2.4 mK
225 K	350 kPa	190 mol/m ³	6.5 mK
300 K	350 kPa	140 mol/m ³	17 mK

Table 1. Best RIGT Monte Carlo quadrature sum thermodynamic temperature uncertainties

For TM11 microwave mode

CAVEATS

- Gas composition impurities, pressure head, thermomolecular pressure correction, pressure gradients along gas flow path, resonator thermal gradients, or any other systematic effects not yet included in calculations ⇒ may increase the overall T_{therm} uncertainty
- δT_{therm} calculated at a single $(p, T, \text{microwave mode})$ combination at a time ⇒ overall T_{therm} uncertainty may be reduced by averaging multiple measurements across modes, along isotherms, etc.
- Special efforts to push p uncertainty below NRC CMCs may reduce overall T_{therm} uncertainty

CONCLUSION

Microwave refractive index gas thermometry uncertainties within the current NRC approach, while not competitive with those of AGT [3], are likely to be small enough between 25 K – 150 K to allow useful measurements of thermodynamic temperature in this regime, to complement those made by AGT and DCGT.

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