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EXPERIMENTAL TESTS OF THE OPTICAL BLOCH EQUATIONS FOR SOLIDS:  
SPECTRAL DIFFUSION

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INTRODUCTION AND REVIEW

Recent free induction decay (FID) studies by DeVoe and Brewer [1] have shown that the conventional optical Bloch equations (OBE) fail to describe the saturation behaviour of the  $^1D_2$  line (592.5 nm) in  $Pr^{3+}:LaF_3$ . This is a surprising and important result both for technological and scientific reasons. For hole-burning memories [2], the effects of such power broadening will play a decisive role in the design of such systems. Scientifically, the study of dephasing (homogeneous linewidth) mechanisms in nuclear and electron magnetic resonance has been actively pursued for some 40 years. In particular, various models of host spin flip induced dephasing have been studied and recently extended to optical transitions [3]. Also several theoretical studies using various models of spin flip induced frequency fluctuations have appeared [4]-[12]. As discussed by Berman [10] however, none of the theories consistently describe the data which includes FID as well as photon echo observations.

Recently we have extended these high resolution FID studies to ruby [13], using an ultra-stable (~ 1 kHz linewidth) dye laser, to subject the theories to further tests with experiment. It might be noted that one contribution to the dephasing spin interaction in ruby occurs between electronic ( $Cr^{3+}$ ) and nuclear ( $Al^{3+}$ ) spins whereas in  $Pr^{3+}:LaF_3$  only nuclear spins are involved.

In addition studies of the power dependence of rotary echoes [14] and hole-burning [15] in the frequency domain have been reported. These results of some of these studies are summarized in Fig. 1-2. The FID ruby data in Fig. 1 are similar to that obtained for  $Pr^{3+}:LaF_3$  in that approximate fits of the Gauss-Markov (GM) [8] and random telegraph (RT) [9] modified OBE could be obtained using a correlation time  $\tau_c=T_2$ , the dephasing time. On the other hand the rotary echo decay was found to disagree with the GM and RT models and was described by the standard OBE. Finally the hole-burning results in Fig. 2, obtained in the frequency domain, are inconsistent with the FID time domain derived hole-widths. However recent theoretical work [16], that assumes a slow, strong stochastic modulation of the ion frequency appears to consistently describe the FID and rotary echo results. Regarding the frequency domain hole studies, preliminary results are presented in

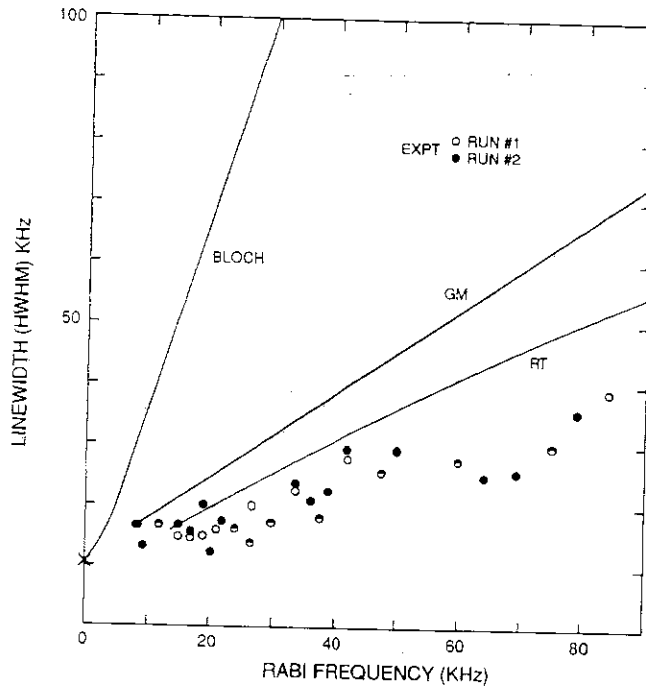


Fig. 1. Optical linewidth (HWHM) of the prepared hole vs Rabi frequency [13]. Theoretical curves for the normal OBE, GM and RT models are shown assuming  $\tau_c = T_2 = 15 \mu\text{sec}$ . The point marked X at zero Rabi frequency was obtained from a photon echo measurement.

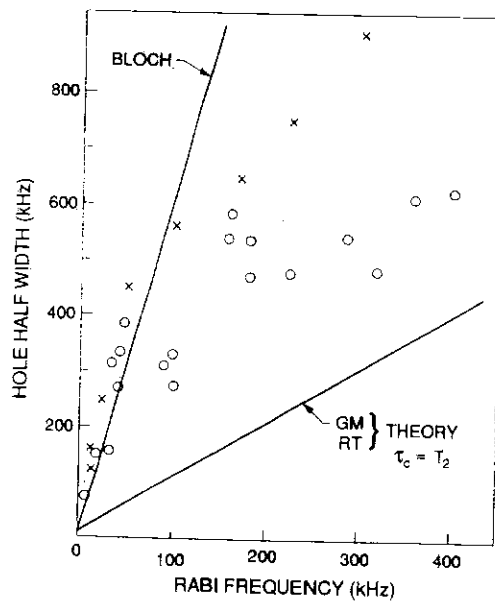


Fig. 2. Optical linewidth of the prepared hole vs Rabi frequency for ruby in the frequency domain [15]. Theoretical curves for the normal OBE, GM and RT models assume  $\tau_c = T_2 = 15 \mu\text{sec}$ . O=Stark shift studies, X=AO frequency shifting.

this paper which show that the earlier [15] hole-widths were broadened by spectral diffusion.

### SPECTRAL DIFFUSION

The apparatus used is shown in Fig. 3. The  ${}^4A_2(-1/2) \rightarrow \bar{E}(-1/2)$  transition [selected by a circular polarizer (CP)] in 0.0034 Wt%  $\text{Cr}_2\text{O}_3$  dilute ruby was studied, as in earlier work [15], with a field of 3.6 kG applied along the c axis and a sample temperature of 2K. A Coherent model 699-21 dye laser was modified [13] to give a peak to peak linewidth  $< 2$  kHz. A pump-probe sequence of laser pulses was used at 25 Hz, with a probe intensity  $\sim 0.1\%$  of the pump and precise probe frequency shifting obtained by a computer-controlled frequency synthesizer driving an acousto-optic (AO) modulator. Typically the pump and probe pulse widths were equal and the boxcar gate width was one-half the pulse width and set on the last half of the probe pulse. The beams before and after the sample were detected and subtracted to reduce amplitude noise and prevent overloading of the amplifier by the pump pulse. Not shown in Fig. 3 are AO switches in front of the diodes to further suppress the pump pulses. Each scan consisted of 101 points taken at 3 second intervals.

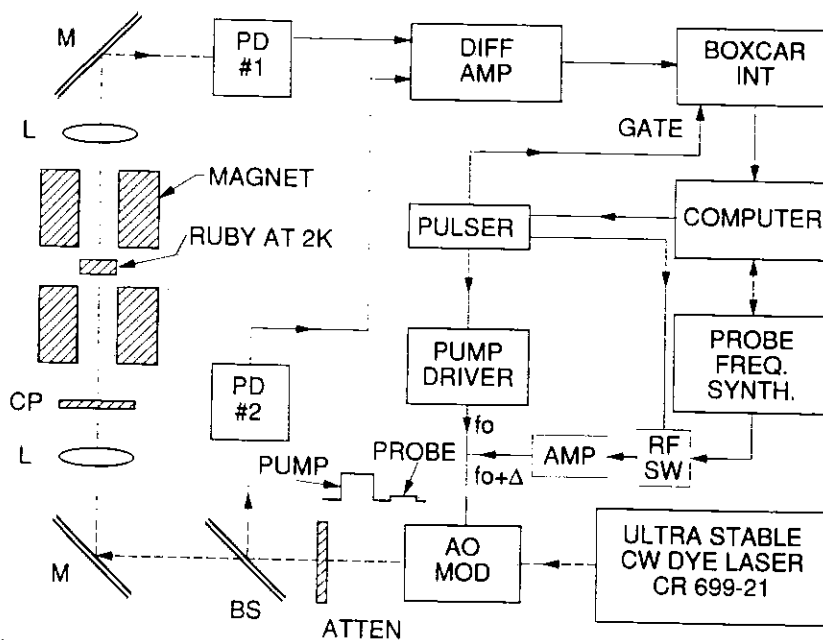


Fig. 3. Apparatus for pump-probe hole-burning studies in the frequency domain (see text).

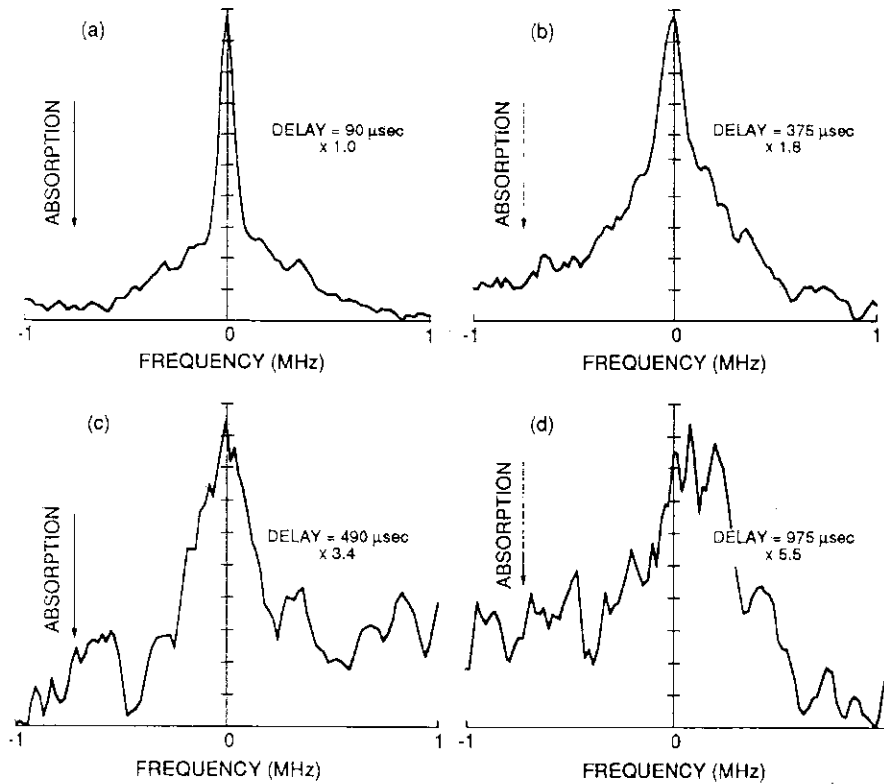


Fig. 4. Hole line shapes observed at various pump-probe delays as defined in the text. The Rabi frequency was 36 kHz and the vertical gain is indicated.

Lineshapes obtained for a series of pump-probe delays (defined as the time from the leading edge of the pump-pulse to the centre of the gate pulse) are shown in Fig. 4. These traces show clear evidence of spectral diffusion over a time scale of  $\sim 300 \mu\text{sec}$ , contrary to the earlier studies of Endo et al [17]. (It is possible that the stated [17] power broadening along with frequency jitter may have obscured the diffusion). These and other data were fit to a Lorentzian line shape by regression analysis. For short times (Fig. 4(a)), only the sharp central peak was used for analysis.

The relation between the observed half-width (HW)  $\Delta\nu_0$  and  $\Delta\nu_h$ , the actual hole HW is

$$\Delta\nu_0 = \Delta\nu_h + \Delta\nu_2 + 2\Delta\nu_p \quad (1)$$

where  $\Delta\nu_2$  is the dephasing HW (12 kHz) and  $\Delta\nu_p$  is the spectral HW due to the finite pulse length. Fig. 5 summarizes the dependence of  $\Delta\nu_h$  on pump-probe delay. Fig. 6 shows, at higher signal to noise ratio, that the lineshape is Lorentzian. A final point on the data is that the diffusion appears to occur similar to that observed [18] for phonon-assisted diffusion where a central sharp spike decays into a broad line (fig. 4(a), (b)). For short delays, it is clear from Fig. 5 that the frequency domain holewidth approaches that inferred from FID [13]. Further studies are planned for various Rabi frequencies to confirm this conclusion. Also studies at high field (to 5T) are planned to clarify the field dependent dephasing observed earlier [19].

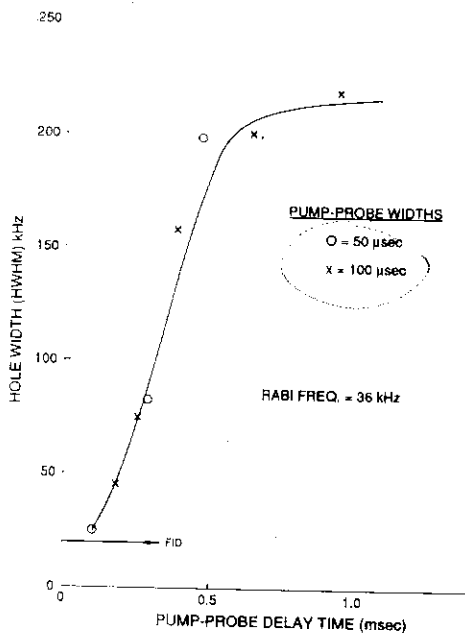


Fig. 5. Plot of deconvoluted hole-width  $\Delta\nu_h$  vs pump-probe delay. The horizontal line indicates the value obtained by FID.

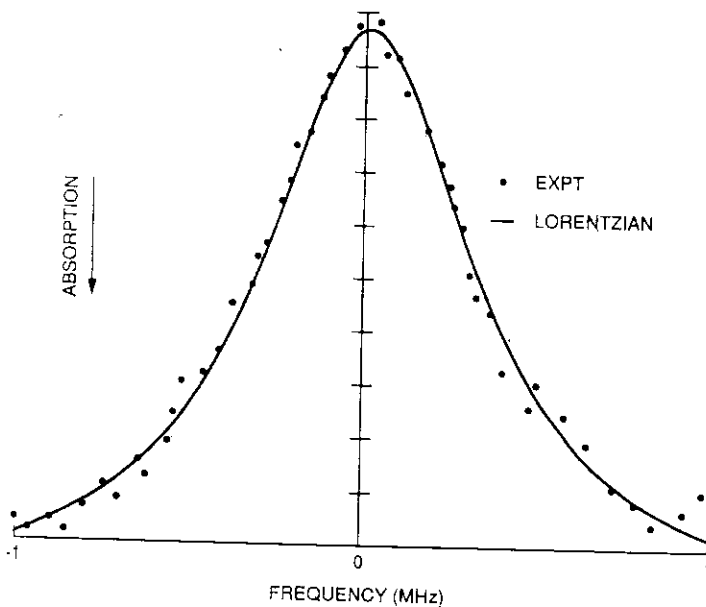


Fig. 6. Regression analysis Lorentzian fit to hole lineshape for a pump-probe delay = 1575  $\mu$ sec. Rabi frequency = 53 kHz. For clarity not all experimental points are plotted.

Concerning the diffusion mechanism, I conclude that  $Al^{27}$  nuclear spin flip-flops outside the frozen core [20] are responsible.  $Cr^{3+}$ - $Cr^{3+}$  spin flips cannot contribute since the flip time for 0.0034% ruby has been measured to be  $\sim 5$  msec [21].

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